

Technical and conceptual advances of LQG

Jerzy Lewandowski
Uniwersytet Warszawski

Challenges for QG, Rome, September 8-12, 2014

Three messages and one anecdote

Three messages and one anecdote

- Quantum geometry coupled to the scalar field: a new operator

$$\hat{q}^{ab} \hat{\phi}_{,a} \hat{\phi}_{,b} \det \hat{q}$$

a new application for LQG, neat, "cute", L, Sahlmann - 2014

Three messages and one anecdote

- Quantum geometry coupled to the scalar field: a new operator

$$\hat{q}^{ab} \hat{\phi}_{,a} \hat{\phi}_{,b} \det \hat{q}$$

a new application for LQG, neat, "cute", L, Sahlmann - 2014

- What is quantum solution to quantum scalar constraint?

Finally we have self adjoint scalar constraint operators

$$\hat{C}(N)$$

themselves. Key for the Dirac quantization of GR L, S - 2014 (coming), Alesci, Assanioussi, L - 2013, A, A, Maekinen, L - 2014 (coming)

Three messages and one anecdote

- Quantum geometry coupled to the scalar field: a new operator

$$\hat{q}^{ab} \hat{\phi}_{,a} \hat{\phi}_{,b} \det \hat{q}$$

a new application for LQG, neat, "cute", L, Sahlmann - 2014

- What is quantum solution to quantum scalar constraint?

Finally we have self adjoint scalar constraint operators

$$\hat{C}(N)$$

themselves. Key for the Dirac quantization of GR L, S - 2014 (coming), Alesci, Assanioussi, L - 2013, A, A, Maekinen, L - 2014 (coming)

- Observables related to distances and angles - the first step - profound, fundamental for general relativistic theories, Duch, Kaminski, L, Świeżewski - 2013, Bodendorfer, D, K, L, Ś - 2014 (coming)

Three messages and one anecdote

- Quantum geometry coupled to the scalar field: a new operator

$$\hat{q}^{ab} \hat{\phi}_{,a} \hat{\phi}_{,b} \det \hat{q}$$

a new application for LQG, neat, "cute", L, Sahlmann - 2014

- What is quantum solution to quantum scalar constraint?

Finally we have self adjoint scalar constraint operators

$$\hat{C}(N)$$

themselves. Key for the Dirac quantization of GR L, S - 2014 (coming), Alesci, Assanioussi, L - 2013, A, A, Maekinen, L - 2014 (coming)

- Observables related to distances and angles - the first step - profound, fundamental for general relativistic theories, Duch, Kaminski, L, Świeżewski - 2013, Bodendorfer, D, K, L, Ś - 2014 (coming)
- A witty anecdote about the maid and gardener - instructive, funny Kurt Vonnegut, Breakfast of Champions - 1973

The connection variables of LQG

The connection variables of LQG

Σ - an underlying 3d manifold, $x^a = x^1, x^2, x^3$ local coordinates

The connection variables of LQG

Σ - an underlying 3d manifold, $x^a = x^1, x^2, x^3$ local coordinates

- The variables:

The connection variables of LQG

Σ - an underlying 3d manifold, $x^a = x^1, x^2, x^3$ local coordinates

- The variables:
 - $\mathfrak{su}(2)$ - the Lie algebra of $SU(2)$, $\tau_i = \tau_1, \tau_2, \tau_3 \in \mathfrak{su}(2)$

The connection variables of LQG

Σ - an underlying 3d manifold, $x^a = x^1, x^2, x^3$ local coordinates

- The variables:
 - $\mathfrak{su}(2)$ - the Lie algebra of $SU(2)$, $\tau_i = \tau_1, \tau_2, \tau_3 \in \mathfrak{su}(2)$
 - (A_a^i, E_i^a) $a, b = 1, 2, 3, i = 1, 2, 3$ - the field variables

The connection variables of LQG

Σ - an underlying 3d manifold, $x^a = x^1, x^2, x^3$ local coordinates

- The variables:
 - $\mathfrak{su}(2)$ - the Lie algebra of $SU(2)$, $\tau_i = \tau_1, \tau_2, \tau_3 \in \mathfrak{su}(2)$
 - (A_a^i, E_i^a) $a, b = 1, 2, 3, i = 1, 2, 3$ - the field variables
 - $\{A_a^i(x), E_i^b(y)\} = \delta_j^i \delta_a^b \delta(x, y)$

The connection variables of LQG

Σ - an underlying 3d manifold, $x^a = x^1, x^2, x^3$ local coordinates

- The variables:
 - $\mathfrak{su}(2)$ - the Lie algebra of $SU(2)$, $\tau_i = \tau_1, \tau_2, \tau_3 \in \mathfrak{su}(2)$
 - (A_a^i, E_i^a) $a, b = 1, 2, 3, i = 1, 2, 3$ - the field variables
 - $\{A_a^i(x), E_i^b(y)\} = \delta_j^i \delta_a^b \delta(x, y)$
 - $\{A, A\} = 0 = \{E, E\}$

The connection variables of LQG

Σ - an underlying 3d manifold, $x^a = x^1, x^2, x^3$ local coordinates

- The variables:
 - $\mathfrak{su}(2)$ - the Lie algebra of $SU(2)$, $\tau_i = \tau_1, \tau_2, \tau_3 \in \mathfrak{su}(2)$
 - (A_a^i, E_i^a) $a, b = 1, 2, 3, i = 1, 2, 3$ - the field variables
 - $\{A_a^i(x), E_i^b(y)\} = \delta_j^i \delta_a^b \delta(x, y)$
 - $\{A, A\} = 0 = \{E, E\}$

- The relation with intrinsic/extrinsic geometry e_a^i / K_{ab} of Σ

The connection variables of LQG

Σ - an underlying 3d manifold, $x^a = x^1, x^2, x^3$ local coordinates

- The variables:
 - $\mathfrak{su}(2)$ - the Lie algebra of $SU(2)$, $\tau_i = \tau_1, \tau_2, \tau_3 \in \mathfrak{su}(2)$
 - (A_a^i, E_i^a) $a, b = 1, 2, 3, i = 1, 2, 3$ - the field variables
 - $\{A_a^i(x), E_i^b(y)\} = \delta_j^i \delta_a^b \delta(x, y)$
 - $\{A, A\} = 0 = \{E, E\}$

- The relation with intrinsic/extrinsic geometry e_a^i / K_{ab} of Σ
Ashtekar, Barbero, Immirzi

The connection variables of LQG

Σ - an underlying 3d manifold, $x^a = x^1, x^2, x^3$ local coordinates

- The variables:
 - $\mathfrak{su}(2)$ - the Lie algebra of $SU(2)$, $\tau_i = \tau_1, \tau_2, \tau_3 \in \mathfrak{su}(2)$
 - (A_a^i, E_i^a) $a, b = 1, 2, 3, i = 1, 2, 3$ - the field variables
 - $\{A_a^i(x), E_i^b(y)\} = \delta_j^i \delta_a^b \delta(x, y)$
 - $\{A, A\} = 0 = \{E, E\}$

- The relation with intrinsic/extrinsic geometry e_a^i / K_{ab} of Σ
Ashtekar, Barbero, Immirzi

$$-2\text{Tr}(\tau_i \tau_j) = \delta_{ij},$$

The connection variables of LQG

Σ - an underlying 3d manifold, $x^a = x^1, x^2, x^3$ local coordinates

- The variables:
 - $\mathfrak{su}(2)$ - the Lie algebra of $SU(2)$, $\tau_i = \tau_1, \tau_2, \tau_3 \in \mathfrak{su}(2)$
 - (A_a^i, E_i^a) $a, b = 1, 2, 3, i = 1, 2, 3$ - the field variables
 - $\{A_a^i(x), E_i^b(y)\} = \delta_j^i \delta_a^b \delta(x, y)$
 - $\{A, A\} = 0 = \{E, E\}$
- The relation with intrinsic/extrinsic geometry e_a^i / K_{ab} of Σ
Ashtekar, Barbero, Immirzi

$$-2\text{Tr}(\tau_i \tau_j) = \delta_{ij},$$

$$A_a^i = \Gamma_a^i + \gamma K_a^i, \quad E_i^a = \frac{1}{16\pi G \gamma} e_b^j e_c^k \epsilon^{abc} \epsilon_{ijk}.$$

Other fields, the constraints

Other fields, the constraints

- Other fields

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

$$\{\phi, \phi\} = 0 = \{\pi, \pi\}$$

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

$$\{\phi, \phi\} = 0 = \{\pi, \pi\}$$

- The constraints (1 - class)

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

$$\{\phi, \phi\} = 0 = \{\pi, \pi\}$$

- The constraints (I - class)

- $\int_\Sigma d^3x N(x) C(x) =: C(N) = C_{\text{gr}}(N) + C_{\text{matt}}(N)$

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

$$\{\phi, \phi\} = 0 = \{\pi, \pi\}$$

- The constraints (I - class)

- $\int_\Sigma d^3x N(x) C(x) =: C(N) = C_{\text{gr}}(N) + C_{\text{matt}}(N)$
- $\int_\Sigma d^3x N^a(x) C_a(x) =: C(\vec{N}) = C_{\text{gr}}(\vec{N}) + C_{\text{matt}}(\vec{N})$

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

$$\{\phi, \phi\} = 0 = \{\pi, \pi\}$$

- The constraints (I - class)

- $\int_\Sigma d^3x N(x) C(x) =: C(N) = C_{\text{gr}}(N) + C_{\text{matt}}(N)$
- $\int_\Sigma d^3x N^a(x) C_a(x) =: C(\vec{N}) = C_{\text{gr}}(\vec{N}) + C_{\text{matt}}(\vec{N})$
- $\int_\Sigma d^3x \Lambda^i G_i(x) =: G(\Lambda) = G_{\text{gr}}(\Lambda) + G_{\text{ferm}}(\Lambda)$

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

$$\{\phi, \phi\} = 0 = \{\pi, \pi\}$$

- The constraints (I - class)
 - $\int_\Sigma d^3x N(x) C(x) =: C(N) = C_{\text{gr}}(N) + C_{\text{matt}}(N)$
 - $\int_\Sigma d^3x N^a(x) C_a(x) =: C(\vec{N}) = C_{\text{gr}}(\vec{N}) + C_{\text{matt}}(\vec{N})$
 - $\int_\Sigma d^3x \Lambda^i G_i(x) =: G(\Lambda) = G_{\text{gr}}(\Lambda) + G_{\text{ferm}}(\Lambda)$
 - Other constraints ($G_{\text{YM}}(\Lambda_{\text{YM}})$)

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

$$\{\phi, \phi\} = 0 = \{\pi, \pi\}$$

- The constraints (I - class)
 - $\int_\Sigma d^3x N(x) C(x) =: C(N) = C_{\text{gr}}(N) + C_{\text{matt}}(N)$
 - $\int_\Sigma d^3x N^a(x) C_a(x) =: C(\vec{N}) = C_{\text{gr}}(\vec{N}) + C_{\text{matt}}(\vec{N})$
 - $\int_\Sigma d^3x \Lambda^i G_i(x) =: G(\Lambda) = G_{\text{gr}}(\Lambda) + G_{\text{ferm}}(\Lambda)$
 - Other constraints ($G_{\text{YM}}(\Lambda_{\text{YM}})$)
- The gauge transformations:

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

$$\{\phi, \phi\} = 0 = \{\pi, \pi\}$$

- The constraints (I - class)
 - $\int_\Sigma d^3x N(x) C(x) =: C(N) = C_{\text{gr}}(N) + C_{\text{matt}}(N)$
 - $\int_\Sigma d^3x N^a(x) C_a(x) =: C(\vec{N}) = C_{\text{gr}}(\vec{N}) + C_{\text{matt}}(\vec{N})$
 - $\int_\Sigma d^3x \Lambda^i G_i(x) =: G(\Lambda) = G_{\text{gr}}(\Lambda) + G_{\text{ferm}}(\Lambda)$
 - Other constraints ($G_{\text{YM}}(\Lambda_{\text{YM}})$)
- The gauge transformations:
 - $\text{Diff}(\Sigma)$

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

$$\{\phi, \phi\} = 0 = \{\pi, \pi\}$$

- The constraints (I - class)
 - $\int_\Sigma d^3x N(x) C(x) =: C(N) = C_{\text{gr}}(N) + C_{\text{matt}}(N)$
 - $\int_\Sigma d^3x N^a(x) C_a(x) =: C(\vec{N}) = C_{\text{gr}}(\vec{N}) + C_{\text{matt}}(\vec{N})$
 - $\int_\Sigma d^3x \Lambda^i G_i(x) =: G(\Lambda) = G_{\text{gr}}(\Lambda) + G_{\text{ferm}}(\Lambda)$
 - Other constraints ($G_{\text{YM}}(\Lambda_{\text{YM}})$)
- The gauge transformations:
 - $\text{Diff}(\Sigma)$
 - $\text{Diff}(\mathbb{R} \times \Sigma)$ orthogonal to (t, Σ)

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

$$\{\phi, \phi\} = 0 = \{\pi, \pi\}$$

- The constraints (I - class)
 - $\int_\Sigma d^3x N(x) C(x) =: C(N) = C_{\text{gr}}(N) + C_{\text{matt}}(N)$
 - $\int_\Sigma d^3x N^a(x) C_a(x) =: C(\vec{N}) = C_{\text{gr}}(\vec{N}) + C_{\text{matt}}(\vec{N})$
 - $\int_\Sigma d^3x \Lambda^i G_i(x) =: G(\Lambda) = G_{\text{gr}}(\Lambda) + G_{\text{ferm}}(\Lambda)$
 - Other constraints ($G_{\text{YM}}(\Lambda_{\text{YM}})$)
- The gauge transformations:
 - $\text{Diff}(\Sigma)$
 - $\text{Diff}(\mathbb{R} \times \Sigma)$ orthogonal to (t, Σ)
 - the frame rotations

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

$$\{\phi, \phi\} = 0 = \{\pi, \pi\}$$

- The constraints (I - class)
 - $\int_\Sigma d^3x N(x) C(x) =: C(N) = C_{\text{gr}}(N) + C_{\text{matt}}(N)$
 - $\int_\Sigma d^3x N^a(x) C_a(x) =: C(\vec{N}) = C_{\text{gr}}(\vec{N}) + C_{\text{matt}}(\vec{N})$
 - $\int_\Sigma d^3x \Lambda^i G_i(x) =: G(\Lambda) = G_{\text{gr}}(\Lambda) + G_{\text{ferm}}(\Lambda)$
 - Other constraints ($G_{\text{YM}}(\Lambda_{\text{YM}})$)
- The gauge transformations:
 - $\text{Diff}(\Sigma)$
 - $\text{Diff}(\mathbb{R} \times \Sigma)$ orthogonal to (t, Σ)
 - the frame rotations
 - Other internal dof rotations (the YM gauge transformations)

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

$$\{\phi, \phi\} = 0 = \{\pi, \pi\}$$

- The constraints (I - class)
 - $\int_\Sigma d^3x N(x) C(x) =: C(N) = C_{\text{gr}}(N) + C_{\text{matt}}(N)$
 - $\int_\Sigma d^3x N^a(x) C_a(x) =: C(\vec{N}) = C_{\text{gr}}(\vec{N}) + C_{\text{matt}}(\vec{N})$
 - $\int_\Sigma d^3x \Lambda^i G_i(x) =: G(\Lambda) = G_{\text{gr}}(\Lambda) + G_{\text{ferm}}(\Lambda)$
 - Other constraints ($G_{\text{YM}}(\Lambda_{\text{YM}})$)
- The gauge transformations:
 - $\text{Diff}(\Sigma)$
 - $\text{Diff}(\mathbb{R} \times \Sigma)$ orthogonal to (t, Σ)
 - the frame rotations
 - Other internal dof rotations (the YM gauge transformations)

The free functions:

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

$$\{\phi, \phi\} = 0 = \{\pi, \pi\}$$

- The constraints (I - class)
 - $\int_\Sigma d^3x N(x) C(x) =: C(N) = C_{\text{gr}}(N) + C_{\text{matt}}(N)$
 - $\int_\Sigma d^3x N^a(x) C_a(x) =: C(\vec{N}) = C_{\text{gr}}(\vec{N}) + C_{\text{matt}}(\vec{N})$
 - $\int_\Sigma d^3x \Lambda^i G_i(x) =: G(\Lambda) = G_{\text{gr}}(\Lambda) + G_{\text{ferm}}(\Lambda)$
 - Other constraints ($G_{\text{YM}}(\Lambda_{\text{YM}})$)
- The gauge transformations:
 - $\text{Diff}(\Sigma)$
 - $\text{Diff}(\mathbb{R} \times \Sigma)$ orthogonal to (t, Σ)
 - the frame rotations
 - Other internal dof rotations (the YM gauge transformations)

The free functions:

- $N \in C(\Sigma)$ - a laps function

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

$$\{\phi, \phi\} = 0 = \{\pi, \pi\}$$

- The constraints (I - class)
 - $\int_\Sigma d^3x N(x) C(x) =: C(N) = C_{\text{gr}}(N) + C_{\text{matt}}(N)$
 - $\int_\Sigma d^3x N^a(x) C_a(x) =: C(\vec{N}) = C_{\text{gr}}(\vec{N}) + C_{\text{matt}}(\vec{N})$
 - $\int_\Sigma d^3x \Lambda^i G_i(x) =: G(\Lambda) = G_{\text{gr}}(\Lambda) + G_{\text{ferm}}(\Lambda)$
 - Other constraints ($G_{\text{YM}}(\Lambda_{\text{YM}})$)
- The gauge transformations:
 - $\text{Diff}(\Sigma)$
 - $\text{Diff}(\mathbb{R} \times \Sigma)$ orthogonal to (t, Σ)
 - the frame rotations
 - Other internal dof rotations (the YM gauge transformations)

The free functions:

- $N \in C(\Sigma)$ - a laps function
- $\vec{N} \in \Gamma(T(\Sigma))$ - a shift vector field

Other fields, the constraints

- Other fields

$$\{\phi_\alpha(x), \pi^\beta(y)\} = \delta_\alpha^\beta \delta(x, y)$$

$$\{\phi, \phi\} = 0 = \{\pi, \pi\}$$

- The constraints (I - class)
 - $\int_\Sigma d^3x N(x) C(x) =: C(N) = C_{\text{gr}}(N) + C_{\text{matt}}(N)$
 - $\int_\Sigma d^3x N^a(x) C_a(x) =: C(\vec{N}) = C_{\text{gr}}(\vec{N}) + C_{\text{matt}}(\vec{N})$
 - $\int_\Sigma d^3x \Lambda^i G_i(x) =: G(\Lambda) = G_{\text{gr}}(\Lambda) + G_{\text{ferm}}(\Lambda)$
 - Other constraints ($G_{\text{YM}}(\Lambda_{\text{YM}})$)
- The gauge transformations:
 - $\text{Diff}(\Sigma)$
 - $\text{Diff}(\mathbb{R} \times \Sigma)$ orthogonal to (t, Σ)
 - the frame rotations
 - Other internal dof rotations (the YM gauge transformations)

The free functions:

- $N \in C(\Sigma)$ - a laps function
- $\vec{N} \in \Gamma(T(\Sigma))$ - a shift vector field
- $\Lambda \in C(\Sigma_{\text{gr}}(2))$

The holonomy-flux variables

Rovelli, Smolin 1988

The holonomy-flux variables

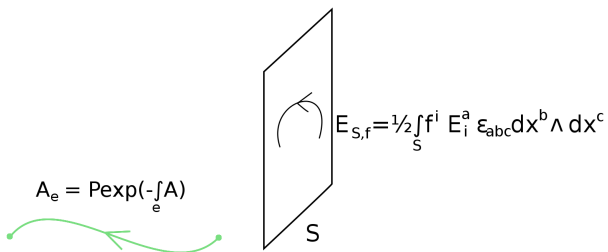
Rovelli, Smolin 1988

$$A_e = P \exp(-\int_e A)$$



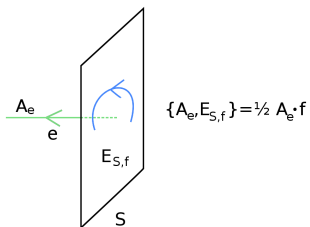
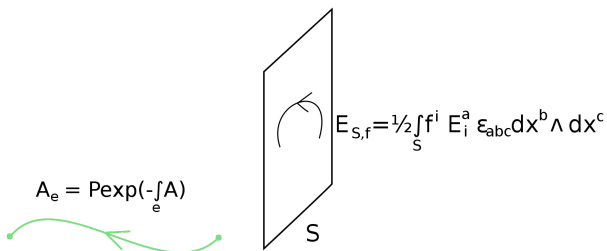
The holonomy-flux variables

Rovelli, Smolin 1988



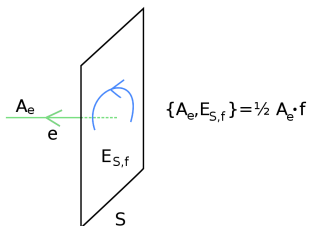
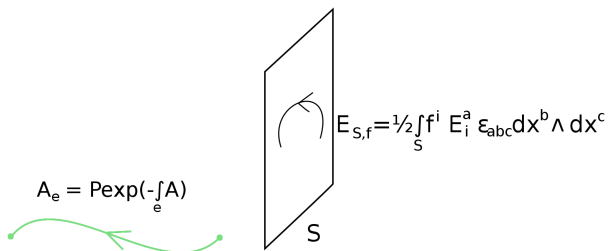
The holonomy-flux variables

Rovelli, Smolin 1988



The holonomy-flux variables

Rovelli, Smolin 1988



e and S ranges curves and, respectively, 2-surfaces in Σ .

The LQG kinematical Hilbert space

Ashtekar, L 1992

The LQG kinematical Hilbert space

Ashtekar, L 1992

$$e : [t_0, t_1] \rightarrow \Sigma$$

The LQG kinematical Hilbert space

Ashtekar, L 1992

$$e : [t_0, t_1] \rightarrow \Sigma$$

$$h_e(A) := \text{Pexp} \int_e -A$$

The LQG kinematical Hilbert space

Ashtekar, L 1992

$$e : [t_0, t_1] \rightarrow \Sigma$$

$$h_e(A) := \text{Pexp} \int_e -A$$

$$\Psi(A) = \psi(h_{e_1}(A), \dots, h_{e_n}(A)), \quad \psi \in \text{Poly}(\text{SU}(2)^n) \quad (1)$$

The LQG kinematical Hilbert space

Ashtekar, L 1992

$$e : [t_0, t_1] \rightarrow \Sigma$$

$$h_e(A) := \text{Pexp} \int_e -A$$

$$\Psi(A) = \psi(h_{e_1}(A), \dots, h_{e_n}(A)), \quad \psi \in \text{Poly}(\text{SU}(2)^n) \quad (1)$$

$\text{Cyl} := \{\Psi \in C(\mathcal{A}) : (1), \{e_1, \dots, e_n\} \text{ embedded graph in } \Sigma\}$

The LQG kinematical Hilbert space

Ashtekar, L 1992

$$e : [t_0, t_1] \rightarrow \Sigma$$

$$h_e(A) := \text{Pexp} \int_e -A$$

$$\Psi(A) = \psi(h_{e_1}(A), \dots, h_{e_n}(A)), \quad \psi \in \text{Poly}(\text{SU}(2)^n) \quad (1)$$

$\text{Cyl} := \{\Psi \in C(\mathcal{A}) : (1), \{e_1, \dots, e_n\} \text{ embedded graph in } \Sigma\}$

$$\int d\mu(A) \Psi(A) = \int dg_1 \dots dg_n \psi(g_1, \dots, g_n)$$

The LQG kinematical Hilbert space

Ashtekar, L 1992

$$e : [t_0, t_1] \rightarrow \Sigma$$

$$h_e(A) := \text{Pexp} \int_e -A$$

$$\Psi(A) = \psi(h_{e_1}(A), \dots, h_{e_n}(A)), \quad \psi \in \text{Poly}(\text{SU}(2)^n) \quad (1)$$

$\text{Cyl} := \{\Psi \in C(\mathcal{A}) : (1), \{e_1, \dots, e_n\} \text{ embedded graph in } \Sigma\}$

$$\int d\mu(A) \Psi(A) = \int dg_1 \dots dg_n \psi(g_1, \dots, g_n)$$

$$\mathcal{H}_{\text{gr}} = L^2(\Omega^{(1)}(\Sigma) \otimes \mathfrak{su}(2), \mu)$$

Spin-networks

Spin-networks

$$U_g \Psi(A) = \Psi(g^{-1}Ag + g^{-1}dg) \quad (2)$$

$$U_g \in \text{Unitary}(\mathcal{H}_{\text{gr}}) \quad (3)$$

Spin-networks

$$U_g \Psi(A) = \Psi(g^{-1}Ag + g^{-1}dg) \quad (2)$$

$$U_g \in \text{Unitary}(\mathcal{H}_{\text{gr}}) \quad (3)$$

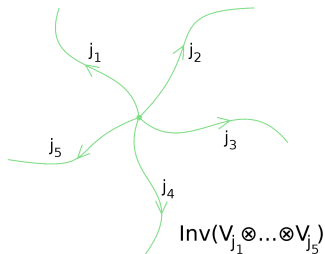
The invariant elements in \mathcal{H}_{gr} *Penrose 1970(?)*, *Rovelli, Smolin*, *Baez 1993*:

Spin-networks

$$U_g \Psi(A) = \Psi(g^{-1}Ag + g^{-1}dg) \quad (2)$$

$$U_g \in \text{Unitary}(\mathcal{H}_{\text{gr}}) \quad (3)$$

The invariant elements in \mathcal{H}_{gr} *Penrose 1970(?)*, *Rovelli, Smolin*,
Baez 1993:

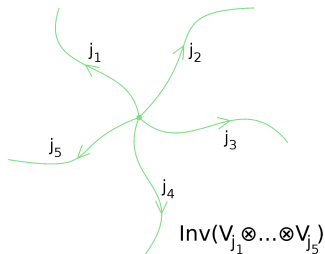


Spin-networks

$$U_g \Psi(A) = \Psi(g^{-1} A g + g^{-1} dg) \quad (2)$$

$$U_g \in \text{Unitary}(\mathcal{H}_{\text{gr}}) \quad (3)$$

The invariant elements in \mathcal{H}_{gr} *Penrose 1970(?)*, *Rovelli, Smolin*,
Baez 1993:



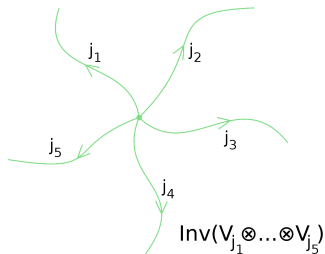
$$\Gamma \ni e_l \mapsto j_l, \quad \Gamma \ni v \mapsto \iota_v \in \text{Inv} V_{j_1} \otimes \dots \otimes V_{j_k}$$

Spin-networks

$$U_g \Psi(A) = \Psi(g^{-1} A g + g^{-1} dg) \quad (2)$$

$$U_g \in \text{Unitary}(\mathcal{H}_{\text{gr}}) \quad (3)$$

The invariant elements in \mathcal{H}_{gr} *Penrose 1970(?)*, *Rovelli, Smolin*,
Baez 1993:



$\Gamma \ni e_l \mapsto j_l$, $\Gamma \ni v \mapsto \iota_v \in \text{Inv} V_{j_1} \otimes \dots \otimes V_{j_k}$
 $|\Gamma, j, \iota\rangle$ - a spin-network state

Scalar field

Scalar field

Classical:

Scalar field

Classical:

$$\{\phi(x), \pi(y)\} = \delta(x, y)$$

Scalar field

Classical:

$$\{\phi(x), \pi(y)\} = \delta(x, y)$$

$$\begin{aligned} C_\phi(x) = & \\ & (8\pi G\gamma)^{-\frac{3}{2}} \frac{\pi^2(x)}{\sqrt{|\det E|}} + \frac{1}{2} (8\pi G\gamma)^{\frac{7}{2}} \phi_{,a} \phi_{,b} E_i^a E_i^b \sqrt{|\det E|} + \\ & (8\pi G\gamma)^{\frac{3}{2}} V(\phi(x)) \sqrt{|\det E|} \end{aligned} \quad (4)$$

Scalar field

Classical:

$$\{\phi(x), \pi(y)\} = \delta(x, y)$$

$$\begin{aligned} C_\phi(x) = & \\ & (8\pi G\gamma)^{-\frac{3}{2}} \frac{\pi^2(x)}{\sqrt{|\det E|}} + \frac{1}{2} (8\pi G\gamma)^{\frac{7}{2}} \phi_{,a} \phi_{,b} E_i^a E_i^b \sqrt{|\det E|} + \\ & (8\pi G\gamma)^{\frac{3}{2}} V(\phi(x)) \sqrt{|\det E|} \end{aligned} \quad (4)$$

Quantum:

Scalar field

Classical:

$$\{\phi(x), \pi(y)\} = \delta(x, y)$$

$$\begin{aligned} C_\phi(x) = & \\ & (8\pi G\gamma)^{-\frac{3}{2}} \frac{\pi^2(x)}{\sqrt{|\det E|}} + \frac{1}{2} (8\pi G\gamma)^{\frac{7}{2}} \phi_{,a} \phi_{,b} E_i^a E_i^b \sqrt{|\det E|} + \\ & (8\pi G\gamma)^{\frac{3}{2}} V(\phi(x)) \sqrt{|\det E|} \end{aligned} \quad (4)$$

Quantum:

$$|\varphi\rangle : \pi \mapsto e^{-\frac{i}{\hbar} \int \pi \varphi} \quad (5)$$

$$\langle \varphi | \varphi' \rangle = \delta_{\varphi\varphi'} \quad (6)$$

$$\hat{\phi}(x)|\varphi\rangle = i\hbar \frac{\delta}{\delta\pi(x)} |\varphi\rangle = \varphi(x)|\varphi\rangle, \quad (7)$$

$$e^{-\frac{i}{\hbar} \int \hat{\pi} \varphi} |\varphi'\rangle = |\varphi + \varphi'\rangle \quad (8)$$

The new operator

L,Sahlmann 2014 (coming)

The new operator

L,Sahlmann 2014 (coming)

$$\mathcal{H}_\phi \otimes \mathcal{H}_{\text{gr}}$$

The new operator

L,Sahlmann 2014 (coming)

$$\mathcal{H}_\phi \otimes \mathcal{H}_{\text{gr}} \ni |\varphi\rangle \otimes |\Gamma, j, \iota\rangle$$

The new operator

L,Sahlmann 2014 (coming)

$$\mathcal{H}_\phi \otimes \mathcal{H}_{\text{gr}} \ni |\varphi\rangle \otimes |\Gamma, j, \iota\rangle$$

$$\int d^3x N(x) \sqrt{\hat{\phi}_{,a} \hat{\phi}_{,b} \hat{q}^{ab} \det \hat{q}} |\varphi\rangle \otimes |\Gamma, j, \iota\rangle =$$

The new operator

L,Sahlmann 2014 (coming)

$$\mathcal{H}_\phi \otimes \mathcal{H}_{\text{gr}} \ni |\varphi\rangle \otimes |\Gamma, j, \iota\rangle$$

$$\begin{aligned} & \int d^3x N(x) \sqrt{\hat{\phi}_{,a} \hat{\phi}_{,b} \hat{q}^{ab} \det \hat{q}} |\varphi\rangle \otimes |\Gamma, j, \iota\rangle = \\ & = (8\pi G\gamma)^2 \left(\sum_I \sqrt{j_I(j_I + 1)} \int_{e_I} N |d\varphi| \right) |\varphi\rangle \otimes |\Gamma, j, \iota\rangle \end{aligned}$$

The new operator

L,Sahlmann 2014 (coming)

$$\mathcal{H}_\phi \otimes \mathcal{H}_{\text{gr}} \ni |\varphi\rangle \otimes |\Gamma, j, \iota\rangle$$

$$\begin{aligned} & \int d^3x N(x) \sqrt{\hat{\phi}_{,a} \hat{\phi}_{,b} \hat{q}^{ab} \det \hat{q}} |\varphi\rangle \otimes |\Gamma, j, \iota\rangle = \\ & = (8\pi G\gamma)^2 \left(\sum_I \sqrt{j_I(j_I + 1)} \int_{e_I} N |d\varphi| \right) |\varphi\rangle \otimes |\Gamma, j, \iota\rangle \end{aligned}$$

Similarity with the QM on a graphs [Rovelli, Vidotto 2008](#) and with the Ma operator [Ma 2002](#)

How to use the very operator

How to use the very operator

We apply it to (set $8\pi G\gamma = 1$)

How to use the very operator

We apply it to (set $8\pi G\gamma = 1$)

$$\hat{\pi}(x) = \pm \sqrt{-\hat{\phi}_{,a}\hat{\phi}_{,b}\hat{E}_i^a\hat{E}_i^b - 2V(\phi(\hat{x}))|\det\hat{E}| - 2\sqrt{|\det\hat{E}|}\hat{C}_{\text{gr}}}$$

How to use the very operator

We apply it to (set $8\pi G\gamma = 1$)

$$\hat{\pi}(x) = \pm \sqrt{-\hat{\phi}_{,a}\hat{\phi}_{,b}\hat{E}_i^a\hat{E}_i^b - 2V(\phi(\hat{x}))|\det\hat{E}|} - 2\sqrt{|\det\hat{E}|\hat{C}_{\text{gr}}}$$

Where, given operators

$$\int_{\Sigma} d^3x \sqrt{\hat{A}(x)}, \quad \int_{\Sigma} d^3x \sqrt{\hat{B}(x)}, \quad \dots$$

How to use the very operator

We apply it to (set $8\pi G\gamma = 1$)

$$\hat{\pi}(x) = \pm \sqrt{-\hat{\phi}_{,a}\hat{\phi}_{,b}\hat{E}_i^a\hat{E}_i^b - 2V(\phi(\hat{x}))|\det\hat{E}|} - 2\sqrt{|\det\hat{E}|\hat{C}_{\text{gr}}}$$

Where, given operators

$$\int_{\Sigma} d^3x \sqrt{\hat{A}(x)}, \quad \int_{\Sigma} d^3x \sqrt{\hat{B}(x)}, \quad \dots$$

We define

How to use the very operator

We apply it to (set $8\pi G\gamma = 1$)

$$\hat{\pi}(x) = \pm \sqrt{-\hat{\phi}_{,a}\hat{\phi}_{,b}\hat{E}_i^a\hat{E}_i^b - 2V(\phi(\hat{x}))|\det\hat{E}|} - 2\sqrt{|\det\hat{E}|\hat{C}_{\text{gr}}}$$

Where, given operators

$$\int_{\Sigma} d^3x \sqrt{\hat{A}(x)}, \quad \int_{\Sigma} d^3x \sqrt{\hat{B}(x)}, \quad \dots$$

We define

$$\int_{\Sigma} d^3x \sqrt{\hat{A}(x) + \hat{B}(x) + \dots} |\varphi\rangle \otimes |\Gamma, j, \iota\rangle :=$$
$$\sum_k \sqrt{\left(\int_{\Sigma_k} d^3x \sqrt{\hat{A}(x)}\right)^2 + \left(\int_{\Sigma_k} d^3x \sqrt{\hat{B}(x)}\right)^2 + \dots} |\varphi\rangle \otimes |\Gamma, j, \iota\rangle \quad (9)$$

How to use the very operator

We apply it to (set $8\pi G\gamma = 1$)

$$\hat{\pi}(x) = \pm \sqrt{-\hat{\phi}_{,a}\hat{\phi}_{,b}\hat{E}_i^a\hat{E}_i^b - 2V(\phi(\hat{x}))|\det\hat{E}| - 2\sqrt{|\det\hat{E}|}\hat{C}_{\text{gr}}}$$

Where, given operators

$$\int_{\Sigma} d^3x \sqrt{\hat{A}(x)}, \quad \int_{\Sigma} d^3x \sqrt{\hat{B}(x)}, \quad \dots$$

We define

$$\int_{\Sigma} d^3x \sqrt{\hat{A}(x) + \hat{B}(x) + \dots} |\varphi\rangle \otimes |\Gamma, j, \iota\rangle :=$$
$$\sum_k \sqrt{\left(\int_{\Sigma_k} d^3x \sqrt{\hat{A}(x)}\right)^2 + \left(\int_{\Sigma_k} d^3x \sqrt{\hat{B}(x)}\right)^2 + \dots} |\varphi\rangle \otimes |\Gamma, j, \iota\rangle \quad (9)$$

provided the RHS is independent of refinements of the partition

$$\Sigma = \bigcup_k \Sigma_k.$$

The miracles of LQG

- The Hilbert space of states of quantum geometry

$$\mathcal{H}_{\text{kin}}$$

with a Diff invariant $(\cdot | \cdot)$

- ∞ -free derivations of operators of quantum geometry

$$\mathcal{O}(\hat{E})$$

- local frame rotation invariant states:

$$\Psi(A) = \Psi(g^{-1}Ag + g^{-1}dg)$$

- Averaging with respect to Diff, and the Hilbert space of Diff invariant states,

$$\mathcal{H}_{\text{Diff}}$$

- ∞ -free derivations of operators of quantum connection

$$\mathcal{O}(\hat{F})$$

in $\mathcal{H}_{\text{Diff}}$ [Rovelli, Smolin 1994](#), [Ashtekar, L 1995](#), [Thiemann 1998](#)

The issue of the $\hat{C}_{\text{gr}}(N)$

The issue of the $\hat{C}_{\text{gr}}(N)$

- $\hat{C}_{\text{gr}}(N)$ not defined in \mathcal{H}_{kin}

The issue of the $\hat{C}_{\text{gr}}(N)$

- $\hat{C}_{\text{gr}}(N)$ not defined in \mathcal{H}_{kin}
- $\hat{C}(N) : \mathcal{H}_{\text{Diff}} \rightarrow$ **suitable dual space**

The issue of the $\hat{C}_{\text{gr}}(N)$

- $\hat{C}_{\text{gr}}(N)$ not defined in \mathcal{H}_{kin}
- $\hat{C}(N) : \mathcal{H}_{\text{Diff}} \rightarrow$ **suitable dual space** , N breaks Diffs

The issue of the $\hat{C}_{\text{gr}}(N)$

- $\hat{C}_{\text{gr}}(N)$ not defined in \mathcal{H}_{kin}
- $\hat{C}(N) : \mathcal{H}_{\text{Diff}} \rightarrow$ suitable dual space , N breaks Diffs
- In principle we can write

$$\hat{C}_{\text{gr}}(N)\psi = 0$$

however:

The issue of the $\hat{C}_{\text{gr}}(N)$

- $\hat{C}_{\text{gr}}(N)$ not defined in \mathcal{H}_{kin}
- $\hat{C}(N) : \mathcal{H}_{\text{Diff}} \rightarrow$ suitable dual space , N breaks Diffs
- In principle we can write

$$\hat{C}_{\text{gr}}(N)\psi = 0$$

however:

- solutions non-normalizable

The issue of the $\hat{C}_{\text{gr}}(N)$

- $\hat{C}_{\text{gr}}(N)$ not defined in \mathcal{H}_{kin}
- $\hat{C}(N) : \mathcal{H}_{\text{Diff}} \rightarrow$ suitable dual space , N breaks Diffs
- In principle we can write

$$\hat{C}_{\text{gr}}(N)\Psi = 0$$

however:

- solutions non-normalizable
- to large space, which Ψ to select?

The issue of the $\hat{C}_{\text{gr}}(N)$

- $\hat{C}_{\text{gr}}(N)$ not defined in \mathcal{H}_{kin}
- $\hat{C}(N) : \mathcal{H}_{\text{Diff}} \rightarrow$ suitable dual space , N breaks Diffs
- In principle we can write

$$\hat{C}_{\text{gr}}(N)\Psi = 0$$

however:

- solutions non-normalizable
- to large space, which Ψ to select?
- $(\Psi|\Psi')_{\text{phys}} = ?$

The issue of the $\hat{C}_{\text{gr}}(N)$

- $\hat{C}_{\text{gr}}(N)$ not defined in \mathcal{H}_{kin}
- $\hat{C}(N) : \mathcal{H}_{\text{Diff}} \rightarrow$ suitable dual space , N breaks Diffs
- In principle we can write

$$\hat{C}_{\text{gr}}(N)\Psi = 0$$

however:

- solutions non-normalizable
- to large space, which Ψ to select?
- $(\Psi|\Psi')_{\text{phys}} = ?$
- For GR coupled to a scalar field we need the Rovelli-Smolín

$$\int_{\Sigma} d^3x \sqrt{-2\sqrt{\det \hat{E}(x)} \hat{C}_{\text{gr}}(x)}$$

The issue of the $\hat{C}_{\text{gr}}(N)$

- $\hat{C}_{\text{gr}}(N)$ not defined in \mathcal{H}_{kin}
- $\hat{C}(N) : \mathcal{H}_{\text{Diff}} \rightarrow$ suitable dual space , N breaks Diffs
- In principle we can write

$$\hat{C}_{\text{gr}}(N)\Psi = 0$$

however:

- solutions non-normalizable
- to large space, which Ψ to select?
- $(\Psi|\Psi')_{\text{phys}} = ?$
- For GR coupled to a scalar field we need the Rovelli-Smolín

$$\int_{\Sigma} d^3x \sqrt{-2\sqrt{\det \hat{E}(x)} \hat{C}_{\text{gr}}(x)}$$

how to define the $\sqrt{\hat{C}_{\text{gr}}(x)}$?

The issue of the $\hat{C}_{\text{gr}}(N)$

- $\hat{C}_{\text{gr}}(N)$ not defined in \mathcal{H}_{kin}
- $\hat{C}(N) : \mathcal{H}_{\text{Diff}} \rightarrow$ suitable dual space , N breaks Diffs
- In principle we can write

$$\hat{C}_{\text{gr}}(N)\Psi = 0$$

however:

- solutions non-normalizable
- to large space, which Ψ to select?
- $(\Psi|\Psi')_{\text{phys}} = ?$
- For GR coupled to a scalar field we need the Rovelli-Smolín

$$\int_{\Sigma} d^3x \sqrt{-2\sqrt{\det \hat{E}(x)} \hat{C}_{\text{gr}}(x)}$$

how to define the $\sqrt{\hat{C}_{\text{gr}}(x)}$?

- Perhaps, we can extend $\mathcal{H}_{\text{Diff}} \subset \mathcal{H}'$ to accomodate $\hat{C}(N)$???

The issue of the $\hat{C}_{\text{gr}}(N)$

- $\hat{C}_{\text{gr}}(N)$ not defined in \mathcal{H}_{kin}
- $\hat{C}(N) : \mathcal{H}_{\text{Diff}} \rightarrow$ suitable dual space , N breaks Diffs
- In principle we can write

$$\hat{C}_{\text{gr}}(N)\Psi = 0$$

however:

- solutions non-normalizable
- to large space, which Ψ to select?
- $(\Psi|\Psi')_{\text{phys}} = ?$
- For GR coupled to a scalar field we need the Rovelli-Smolín

$$\int_{\Sigma} d^3x \sqrt{-2\sqrt{\det \hat{E}(x)} \hat{C}_{\text{gr}}(x)}$$

how to define the $\sqrt{\hat{C}_{\text{gr}}(x)}$?

- Perhaps, we can extend $\mathcal{H}_{\text{Diff}} \subset \mathcal{H}'$ to accommodate $\hat{C}(N)$???
- The obstacle:

The issue of the $\hat{C}_{\text{gr}}(N)$

- $\hat{C}_{\text{gr}}(N)$ not defined in \mathcal{H}_{kin}
- $\hat{C}(N) : \mathcal{H}_{\text{Diff}} \rightarrow$ suitable dual space , N breaks Diffs
- In principle we can write

$$\hat{C}_{\text{gr}}(N)\Psi = 0$$

however:

- solutions non-normalizable
- to large space, which Ψ to select?
- $(\Psi|\Psi')_{\text{phys}} = ?$
- For GR coupled to a scalar field we need the Rovelli-Smolín

$$\int_{\Sigma} d^3x \sqrt{-2\sqrt{\det \hat{E}(x)} \hat{C}_{\text{gr}}(x)}$$

how to define the $\sqrt{\hat{C}_{\text{gr}}(x)}$?

- Perhaps, we can extend $\mathcal{H}_{\text{Diff}} \subset \mathcal{H}'$ to accommodate $\hat{C}(N)$???
- The obstacle: Ψ Diff invariant $\Rightarrow (\hat{C}(N)\Psi | \hat{C}(N)\Psi)' = 0$

The issue of the $\hat{C}_{\text{gr}}(N)$

- $\hat{C}_{\text{gr}}(N)$ not defined in \mathcal{H}_{kin}
- $\hat{C}(N) : \mathcal{H}_{\text{Diff}} \rightarrow$ suitable dual space , N breaks Diffs
- In principle we can write

$$\hat{C}_{\text{gr}}(N)\Psi = 0$$

however:

- solutions non-normalizable
- to large space, which Ψ to select?
- $(\Psi|\Psi')_{\text{phys}} = ?$
- For GR coupled to a scalar field we need the Rovelli-Smolín

$$\int_{\Sigma} d^3x \sqrt{-2\sqrt{\det \hat{E}(x)} \hat{C}_{\text{gr}}(x)}$$

how to define the $\sqrt{\hat{C}_{\text{gr}}(x)}$?

- Perhaps, we can extend $\mathcal{H}_{\text{Diff}} \subset \mathcal{H}'$ to accommodate $\hat{C}(N)$???
- The obstacle: Ψ Diff invariant $\Rightarrow (\hat{C}(N)\Psi | \hat{C}(N)\Psi)' = 0$ for every $N \in C_0^\infty(\Sigma)$. L, Marolf 1999

Partial ways out

Partial ways out

- Take $N = 1$,

Partial ways out

- Take $N = 1$,

$$\hat{C}(1) : \mathcal{H}_{\text{Diff}} \rightarrow \mathcal{H}_{\text{Diff}}$$

Partial ways out

- Take $N = 1$,

$$\hat{C}(1) : \mathcal{H}_{\text{Diff}} \rightarrow \mathcal{H}_{\text{Diff}}$$

$$\hat{C}_{\text{gr,sym}} := \frac{1}{2} \left(\hat{C}_{\text{gr}}(1) + \hat{C}_{\text{gr}}^\dagger(1) \right).$$

Partial ways out

- Take $N = 1$,

$$\hat{C}(1) : \mathcal{H}_{\text{Diff}} \rightarrow \mathcal{H}_{\text{Diff}}$$

$$\hat{C}_{\text{gr,sym}} := \frac{1}{2} \left(\hat{C}_{\text{gr}}(1) + \hat{C}_{\text{gr}}^\dagger(1) \right).$$

But $N = 1$ is not enough...

Partial ways out

- Take $N = 1$,

$$\hat{C}(1) : \mathcal{H}_{\text{Diff}} \rightarrow \mathcal{H}_{\text{Diff}}$$

$$\hat{C}_{\text{gr,sym}} := \frac{1}{2} \left(\hat{C}_{\text{gr}}(1) + \hat{C}_{\text{gr}}^\dagger(1) \right).$$

But $N = 1$ is not enough...

- Try to define directly

Partial ways out

- Take $N = 1$,

$$\hat{C}(1) : \mathcal{H}_{\text{Diff}} \rightarrow \mathcal{H}_{\text{Diff}}$$

$$\hat{C}_{\text{gr,sym}} := \frac{1}{2} \left(\hat{C}_{\text{gr}}(1) + \hat{C}_{\text{gr}}^\dagger(1) \right).$$

But $N = 1$ is not enough...

- Try to define directly
 - either

$$\int d^3x \sqrt{-\det \hat{E}(x)} \hat{C}(x)$$

Partial ways out

- Take $N = 1$,

$$\hat{C}(1) : \mathcal{H}_{\text{Diff}} \rightarrow \mathcal{H}_{\text{Diff}}$$

$$\hat{C}_{\text{gr,sym}} := \frac{1}{2} \left(\hat{C}_{\text{gr}}(1) + \hat{C}_{\text{gr}}^\dagger(1) \right).$$

But $N = 1$ is not enough...

- Try to define directly
 - either

$$\int d^3x \sqrt{-\det \hat{E}(x)} \hat{C}(x)$$

- or

Partial ways out

- Take $N = 1$,

$$\hat{C}(1) : \mathcal{H}_{\text{Diff}} \rightarrow \mathcal{H}_{\text{Diff}}$$

$$\hat{C}_{\text{gr,sym}} := \frac{1}{2} \left(\hat{C}_{\text{gr}}(1) + \hat{C}_{\text{gr}}^\dagger(1) \right).$$

But $N = 1$ is not enough...

- Try to define directly
 - either

$$\int d^3x \sqrt{-\det \hat{E}(x)} \hat{C}(x)$$

- or

$$\int d^3x \frac{\hat{C}^2(x)}{\sqrt{\det \hat{E}(x)}}$$

Partial ways out

- Take $N = 1$,

$$\hat{C}(1) : \mathcal{H}_{\text{Diff}} \rightarrow \mathcal{H}_{\text{Diff}}$$

$$\hat{C}_{\text{gr,sym}} := \frac{1}{2} \left(\hat{C}_{\text{gr}}(1) + \hat{C}_{\text{gr}}^\dagger(1) \right).$$

But $N = 1$ is not enough...

- Try to define directly
 - either

$$\int d^3x \sqrt{-\det \hat{E}(x)} \hat{C}(x)$$

- or

$$\int d^3x \frac{\hat{C}^2(x)}{\sqrt{\det \hat{E}(x)}}$$

the Master constraint program (**Thiemann**)

The new Hilbert space

L, Sahlmann 2014 (coming)

The new Hilbert space

L, Sahlmann 2014 (coming)

$$\mathcal{H}_{\text{new}} = \bigoplus_{\{x_1, \dots, x_k\} \subset \Sigma} \mathcal{H}_{\{x_1, \dots, x_k\}}$$

The new Hilbert space

L, Sahlmann 2014 (coming)

$$\mathcal{H}_{\text{new}} = \bigoplus_{\{x_1, \dots, x_k\} \subset \Sigma} \mathcal{H}_{\{x_1, \dots, x_k\}}$$

$\mathcal{H}_{\{x_1, \dots, x_k\}}$ is spanned by all the $|\Gamma, j, \iota\rangle$ based at $\{x_1, \dots, x_k\}$,

The new Hilbert space

L, Sahlmann 2014 (coming)

$$\mathcal{H}_{\text{new}} = \bigoplus_{\{x_1, \dots, x_k\} \subset \Sigma} \mathcal{H}_{\{x_1, \dots, x_k\}}$$

$\mathcal{H}_{\{x_1, \dots, x_k\}}$ is spanned by all the $|\Gamma, j, \iota\rangle$ based at $\{x_1, \dots, x_k\}$, averaged with respect to $\text{Diff}(\Sigma, \{x_1, \dots, x_k\})$.

Operators in \mathcal{H}_{new}

Operators in \mathcal{H}_{new}

We define operators in \mathcal{H}_{new} either by passing them by the duality from the kinematical \mathcal{H}_{kin} ,

Operators in \mathcal{H}_{new}

We define operators in \mathcal{H}_{new} either by passing them by the duality from the kinematical \mathcal{H}_{kin} , or $\hat{C}_{\text{gr}}(N)$ from scratch:

Operators in \mathcal{H}_{new}

We define operators in \mathcal{H}_{new} either by passing them by the duality from the kinematical \mathcal{H}_{kin} , or $\hat{C}_{\text{gr}}(N)$ from scratch:

$$d^3x N(x) \hat{O}(x) = \sum_{x \in \Sigma} N(x) \hat{O}_x$$

Operators in \mathcal{H}_{new}

We define operators in \mathcal{H}_{new} either by passing them by the duality from the kinematical \mathcal{H}_{kin} , or $\hat{C}_{\text{gr}}(N)$ from scratch:

$$d^3x N(x) \hat{O}(x) = \sum_{x \in \Sigma} N(x) \hat{O}_x$$

$$\hat{O}_x : \mathcal{H}_{\{x_1, \dots, x_k\}} \rightarrow \mathcal{H}_{\{x_1, \dots, x_k\}}$$

Operators in \mathcal{H}_{new}

We define operators in \mathcal{H}_{new} either by passing them by the duality from the kinematical \mathcal{H}_{kin} , or $\hat{C}_{\text{gr}}(N)$ from scratch:

$$d^3x N(x) \hat{O}(x) = \sum_{x \in \Sigma} N(x) \hat{O}_x$$

$$\hat{O}_x : \mathcal{H}_{\{x_1, \dots, x_k\}} \rightarrow \mathcal{H}_{\{x_1, \dots, x_k\}}$$

$$\hat{O}_x : \mathcal{H}_{\{x_1, \dots, x_k\}} \rightarrow 0 \text{ unless } x = x_1, \dots, x_k$$

Operators in \mathcal{H}_{new}

We define operators in \mathcal{H}_{new} either by passing them by the duality from the kinematical \mathcal{H}_{kin} , or $\hat{C}_{\text{gr}}(N)$ from scratch:

$$d^3x N(x) \hat{O}(x) = \sum_{x \in \Sigma} N(x) \hat{O}_x$$

$$\hat{O}_x : \mathcal{H}_{\{x_1, \dots, x_k\}} \rightarrow \mathcal{H}_{\{x_1, \dots, x_k\}}$$

$$\hat{O}_x : \mathcal{H}_{\{x_1, \dots, x_k\}} \rightarrow 0 \text{ unless } x = x_1, \dots, x_k$$

$$\mathcal{O}(x) = \sqrt{\det E(x)}, \sqrt{\det E(x)} \text{Ric}(x), C_{\text{gr}}(x)$$

The quantum scalar constraint

With this result we can:

- Consider $\hat{C}^\dagger(N)$ (it has sufficiently large domain).
- Define

$$\hat{C}(x)_{\text{sym}} = \frac{1}{2}(\hat{C} + \hat{C}^\dagger)$$

- Find a self adjoint extension $\hat{C}_{\text{s.a.}}$.
- Spectrally expand each

$$\mathcal{H}_{\{x_1, \dots, x_k\}} = \int^\oplus dc_1 \dots dc_k \mathcal{H}_{\{x_1, \dots, x_k\}}^{c_1 \dots c_k}$$

- promote $\mathcal{H}_{\{x_1, \dots, x_k\}}^{0 \dots 0}$ to be solutions to the quantum scalar constraint

Solutions to the matter free LQG

The elements of $\mathcal{H}_{\{x_1, \dots, x_k\}}$ have to be further averaged with respect to the vertex **non**-preserving $\text{Diff}(\Sigma)$ (the vector constraint),

$$\mathcal{H}_{\{x_1, \dots, x_k\}} \eta_{\text{new}}(\Psi) \mapsto \frac{1}{k} \sum_{[f] \in \text{Diff}(\Sigma)/\text{Diff}(\Sigma)} \eta_{\text{new}}(U_f \Psi)$$

The map passes to the $\mathcal{H}_{\{x_1, \dots, x_k\}}^{0 \dots 0}$ space, and it's image defines a Hilbert space

$$\mathcal{H}_{(k)}$$

of k -vertex solutions to the scalar **AND** vector constraint.
The full Hilbert space is

$$\mathcal{H}_{\text{phys}} = \bigoplus_k \mathcal{H}_{(k)}$$

Observables - emergence of points

Duch, Kaminski, L, Świeżewski 2014, Bodendorfer, L, Świeżewski 2014 (coming)

Observables - emergence of points

Duch, Kaminski, L, Świeżewski 2014, Bodendorfer, L, Świeżewski 2014 (coming)

Gravitational ADM data q_{ab}, p^{ab} and other fields ϕ, π .

Observables - emergence of points

Duch, Kaminski, L, Świeżewski 2014, Bodendorfer, L, Świeżewski 2014 (coming)

Gravitational ADM data q_{ab}, p^{ab} and other fields ϕ, π .

Observer:

Observables - emergence of points

Duch, Kaminski, L, Świeżewski 2014, Bodendorfer, L, Świeżewski 2014 (coming)

Gravitational ADM data q_{ab}, p^{ab} and other fields ϕ, π .

Observer:

- $x_0 \in \Sigma$

Observables - emergence of points

Duch, Kaminski, L, Świążewski 2014, Bodendorfer, L, Świążewski 2014 (coming)

Gravitational ADM data q_{ab}, p^{ab} and other fields ϕ, π .

Observer:

- $x_0 \in \Sigma$
- \mathbb{R}^3 equipped with spherical coordinates r, θ^1, θ^2

Observables - emergence of points

Duch, Kaminski, L, Świążewski 2014, Bodendorfer, L, Świążewski 2014 (coming)

Gravitational ADM data q_{ab}, p^{ab} and other fields ϕ, π .

Observer:

- $x_0 \in \Sigma$
- \mathbb{R}^3 equipped with spherical coordinates r, θ^1, θ^2
- for every metric q in Σ , a map

$$q \mapsto f_q : \mathbb{R}^3 \rightarrow \Sigma,$$

Observables - emergence of points

Duch, Kaminski, L, Świeżewski 2014, Bodendorfer, L, Świeżewski 2014 (coming)

Gravitational ADM data q_{ab}, p^{ab} and other fields ϕ, π .

Observer:

- $x_0 \in \Sigma$
- \mathbb{R}^3 equipped with spherical coordinates r, θ^1, θ^2
- for every metric q in Σ , a map

$$q \mapsto f_q : \mathbb{R}^3 \rightarrow \Sigma,$$

where f such that

Observables - emergence of points

Duch, Kaminski, L, Świążewski 2014, Bodendorfer, L, Świążewski 2014 (coming)

Gravitational ADM data q_{ab}, p^{ab} and other fields ϕ, π .

Observer:

- $x_0 \in \Sigma$
- \mathbb{R}^3 equipped with spherical coordinates r, θ^1, θ^2
- for every metric q in Σ , a map

$$q \mapsto f_q : \mathbb{R}^3 \rightarrow \Sigma,$$

where f such that

$$f_q : (0, 0, 0) \mapsto x_0$$

Observables - emergence of points

Duch, Kaminski, L, Świeżewski 2014, Bodendorfer, L, Świeżewski 2014 (coming)

Gravitational ADM data q_{ab}, p^{ab} and other fields ϕ, π .

Observer:

- $x_0 \in \Sigma$
- \mathbb{R}^3 equipped with spherical coordinates r, θ^1, θ^2
- for every metric q in Σ , a map

$$q \mapsto f_q : \mathbb{R}^3 \rightarrow \Sigma,$$

where f such that

$$f_q : (0, 0, 0) \mapsto x_0$$

f_q : radial geodesic line \rightarrow radial geodesic line

Observables - emergence of points

Duch, Kaminski, L, Świeżewski 2014, Bodendorfer, L, Świeżewski 2014 (coming)

Gravitational ADM data q_{ab}, p^{ab} and other fields ϕ, π .

Observer:

- $x_0 \in \Sigma$
- \mathbb{R}^3 equipped with spherical coordinates r, θ^1, θ^2
- for every metric q in Σ , a map

$$q \mapsto f_q : \mathbb{R}^3 \rightarrow \Sigma,$$

where f such that

$$f_q : (0, 0, 0) \mapsto x_0$$

f_q : radial geodesic line \rightarrow radial geodesic line

r coincides with the distance according to q

Observables - emergence of points

Duch, Kaminski, L, Świeżewski 2014, Bodendorfer, L, Świeżewski 2014 (coming)

Gravitational ADM data q_{ab}, p^{ab} and other fields ϕ, π .

Observer:

- $x_0 \in \Sigma$
- \mathbb{R}^3 equipped with spherical coordinates r, θ^1, θ^2
- for every metric q in Σ , a map

$$q \mapsto f_q : \mathbb{R}^3 \rightarrow \Sigma,$$

where f such that

$$f_q : (0, 0, 0) \mapsto x_0$$

f_q : radial geodesic line \rightarrow radial geodesic line

r coincides with the distance according to q

The physical points are triples (r, θ^1, θ^2)

Observables - the observables

Observables - the observables

Given a scalar field ϕ is Σ ,

Observables - the observables

Given a scalar field ϕ on Σ , the observable is

Observables - the observables

Given a scalar field ϕ on Σ , the observable is

$$\Phi(r, \theta, q, p, \phi, \pi) := \phi(f_q(r, \theta))$$

Observables - the observables

Given a scalar field ϕ is Σ , the observable is

$$\Phi(r, \theta, q, p, \phi, \pi) := \phi(f_q(r, \theta))$$

Given indices $i, j = r, 1, 2$

Observables - the observables

Given a scalar field ϕ on Σ , the observable is

$$\Phi(r, \theta, q, p, \phi, \pi) := \phi(f_q(r, \theta))$$

Given indices $i, j = r, 1, 2$

$$Q_{ij}(r, \theta, q, p, \phi, \pi) := (f^*q)_{ij}(r, \theta).$$

Observables - the observables

Given a scalar field ϕ is Σ , the observable is

$$\Phi(r, \theta, q, p, \phi, \pi) := \phi(f_q(r, \theta))$$

Given indices $i, j = r, 1, 2$

$$Q_{ij}(r, \theta, q, p, \phi, \pi) := (f^* q)_{ij}(r, \theta).$$

Similarly for P^{ij} and Π .

Result

Result

Gauge constraints

$$Q_{ri} = \delta_{ri}$$

Result

Gauge constraints

$$Q_{ri} = \delta_{ri}$$

The variables

Result

Gauge constraints

$$Q_{ri} = \delta_{ri}$$

The variables

$$Q_{AB}(r, \theta, q, p, \phi, \pi), P^{AB}(r, \theta, q, p, \phi, \pi)$$

Result

Gauge constraints

$$Q_{ri} = \delta_{ri}$$

The variables

$$Q_{AB}(r, \theta, q, p, \phi, \pi), P^{AB}(r, \theta, q, p, \phi, \pi)$$

$$\Phi(r, \theta, q, p, \phi, \pi), \Pi(r, \theta, q, p, \phi, \pi)$$

canonically conjugate (be careful about x_0 , regularity).

Result

Gauge constraints

$$Q_{ri} = \delta_{ri}$$

The variables

$$Q_{AB}(r, \theta, q, p, \phi, \pi), P^{AB}(r, \theta, q, p, \phi, \pi)$$

$$\Phi(r, \theta, q, p, \phi, \pi), \Pi(r, \theta, q, p, \phi, \pi)$$

canonically conjugate (be careful about x_0 , regularity).

The remaining vabs

$$P^{ri}(r, \theta, q, p, \phi, \pi)$$

Result

Gauge constraints

$$Q_{ri} = \delta_{ri}$$

The variables

$$Q_{AB}(r, \theta, q, p, \phi, \pi), P^{AB}(r, \theta, q, p, \phi, \pi)$$

$$\Phi(r, \theta, q, p, \phi, \pi), \Pi(r, \theta, q, p, \phi, \pi)$$

canonically conjugate (be careful about x_0 , regularity).

The remaining vabs

$$P^{ri}(r, \theta, q, p, \phi, \pi)$$

determined on the vector constraint surface

$$P^{ri}(r, \theta, q, p, \phi, \pi) = \tilde{P}^{ri}(r, \theta, Q_{AB}, P^{AB}, \Phi, \Pi)$$

Summary, outlook